

## Moments of random states

$$\mathbb{E} \left| \langle x | x \rangle \right|^n = \frac{\text{Tr}_{\text{sym}}}{\text{Tr} P_{\text{sym}}} = \frac{\frac{1}{n!} \sum_{\pi \in S_n} P_{\pi}}{\binom{d+n-1}{n}} = \frac{\sum_{\pi \in S_n} P_{\pi}}{d(d+1) \cdots (d+n-1)}$$

$\text{Tr}_{\text{sym}} = \sum_{\text{types } t} \text{Tr}_t X_t X_t^*$

$$\mathbb{E} \left| \langle y | y \rangle \right|^M = \frac{\sum_{\pi} P_{\pi}}{d^n} \quad (\text{Wick's thm})$$

e.g.  $\mathbb{E} |\langle u | 0 \rangle|^{2k} = \frac{1}{\binom{d+k-1}{k}}$

$k=1$	$\frac{1}{d}$
$k=2$	$\frac{2}{d(d+1)}$

### Application

"anticentration"

$$\mathbb{E} \sum_{\substack{x \in \{0,1\}^n \\ \psi \times e \{0,1\}^n}} |\langle \psi | x \rangle|^4 = \text{Tr} \sum_x \text{Tr}_{x_1 \otimes x_2 \otimes \dots \otimes x_d} \text{Tr}_{\text{sym}} \cdot \frac{2}{2^n(2^n+1)} = \frac{2}{2^n+1} \approx \frac{2}{2^n} \quad \text{vs } \frac{1}{2^n} \text{ uniform}$$

### R\uacute{e}nyi entropies

$$H_{\alpha}(p) = \frac{1}{1-\alpha} \log \sum_x p(x)^{\alpha} \quad S_{\alpha}(p) = \frac{1}{1-\alpha} \log \text{Tr} p^{\alpha}$$

properties

$$\frac{\partial}{\partial \alpha} H_{\alpha} \leq 0 \quad \text{so } H_{\alpha}(\text{unif}_d) = \log d$$

Schur-concave

$$0 \leq H_{\alpha} \leq \log d$$

pure

Id

$\alpha < 1$  sensitive to low values  
 $\alpha > 1$  "high values"

$$S_0 = \log \text{rank } p$$

$$S_1 = S(p)$$

$$S_2 = -\log \text{Tr } p^2$$

$$S_{\infty} = -\log \|p\| \geq \frac{1}{2} S_2(p)$$

$$p(x) = |\langle x | \psi \rangle|^2$$

$$H(p) \geq H_2(p) \approx n-1$$

more precisely

$$\mathbb{E} H_2(x) = \mathbb{E} -\log \sum_x p(x)^2 \geq -\log \mathbb{E} p(x)^2 \approx n-1$$

irrelevant for q. supremacy via random circuit sampling

Similarly, most pure states are very entangled

$$|\psi\rangle \in \mathbb{C}^{d_A} \otimes \mathbb{C}^{d_B} \text{ random } \left( \text{e.g. } d_A = 2^{n_A}, d_B = 2^{n_B} \right)$$

$$\mathbb{E} S(A) \geq \mathbb{E} S_2(A) = -\mathbb{E} \log \text{tr} \Psi_A^2 \geq -\log \mathbb{E} \text{tr} \Psi_A^2$$

relabel trick:  $\text{tr } F(X \otimes Y) = \text{tr } XY$

$$\text{tr} \Psi_A^2 = \text{tr} (\Psi_A \otimes \Psi_A) F = \text{tr} (\Psi_{A_1} \otimes \Psi_{A_2}) (F_{A_1 A_2} \otimes I_{B_1 B_2})$$

$$\begin{aligned} \mathbb{E} \text{tr} \Psi_A^2 &= \text{tr} \mathbb{E} (\Psi_{A_1 B_1} \otimes \Psi_{A_2 B_2}) (F_{A_1 A_2} \otimes I_{B_1 B_2}) \\ &= \text{tr} \underbrace{I + F_{A_1 A_2} F_{B_1 B_2}}_{D(D+1)} \cdot F_{A_1 A_2} = \frac{d_A d_B^2 + d_A^2 d_B}{d_A d_B (d_A d_B + 1)} = \frac{d_A + d_B}{d_A d_B + 1} \approx \frac{1}{d_A} + \frac{1}{d_B} \end{aligned}$$

$$\mathbb{E} S(A) \geq -\log \left( \frac{1}{d_A} + \frac{1}{d_B} \right)$$

$$d_A = d_B \Rightarrow \mathbb{E} S(A) \geq \log d_A - 1$$

$$d_B > d_A \quad \mathbb{E} S(A) \geq \log d_B - \frac{d_A}{d_B}$$

$$\begin{aligned} \left\| \Psi_A - \frac{I}{d_A} \right\|_1^2 &\leq d_A \left\| \Psi_A - \frac{I}{d_A} \right\|_2^2 = d_A \text{tr} \left( \Psi_A - \frac{I}{d_A} \right)^2 \\ &= d_A \text{tr} \Psi_A^2 - 1 \approx \frac{d_A}{d_B} \end{aligned}$$

group representations

$$R : G \rightarrow U(V)$$

$$\text{s.t. } R(gh) = R(g)R(h)$$

a group unitary operators on a vector space  $V$

$$V^G = \text{invariant subspace} = \left\{ |\psi\rangle \in V : R(g)|\psi\rangle = |\psi\rangle \quad \forall g \in G \right\}$$

$$\Pi = \bigcap_{|G|} \sum_{g \in G} R(g)$$

$$\underline{\text{claim}} \quad \Pi = \text{proj } V^G$$

$$R(h)\Pi = \Pi$$

$$\underline{\text{pf}} \quad R(h) \sum_g R(g) = \sum_g R(hg) = \sum_g R(g) \quad \text{since } g \rightarrow hg \text{ is 1-1}$$

(or measure-preserving if

$G$  is a compact group)

$$\underline{\text{pf}} \quad \Pi^\dagger \Pi = |G|^{-1} \sum_{h \in G} R(h)^\dagger R(h) = \Pi \quad \Rightarrow \quad \Pi \text{ is a projector}$$

$$R(h)\Pi|\psi\rangle = \Pi|\psi\rangle \Rightarrow \Pi|\psi\rangle \in V^G$$

$$\text{conversely if } |\psi\rangle \in V^G \text{ then } \Pi|\psi\rangle = \frac{1}{|G|} \sum_g R(g)|\psi\rangle = |\psi\rangle$$