

# Monogamy without symmetry

$$I(A; B_1 \dots B_k) = \sum_{i=1}^k I(A; B_i | B_{\setminus i}) \quad B_{\setminus i} = B_1 \dots B_{i-1} B_{i+1} \dots B_k$$

$$\leq 2 \log d_A$$

$$\Rightarrow \mathbb{E}_i I(A; B_i) \leq \frac{2 \log d_A}{k}$$

"squashed entanglement"  
 $E_{sq}(\rho_{AB}) = \min_E \frac{1}{2} I(A; B | E)_\rho$   
 $\tilde{\rho}_{ABE} : \tilde{\rho}_{AB} = \rho_{AB}$

so what if  $I(A; C | B) \leq \epsilon$  ?  
 what if  $I(A; C | B) = 0$  ?

classically  $\text{CMI} = 0 \Leftrightarrow p(x, y, z) = e^{-F_1(x, y) - F_2(y, z)}$  etc...

quantumly

$$I(A; C | B) = 0 \Leftrightarrow \rho = \exp(\log \rho^{AB} + \log \rho^{BC} - \log \rho^B)$$

$$\Leftrightarrow B \cong \bigoplus_{\alpha} B_{\alpha}^L \otimes B_{\alpha}^R$$

↑ classical variable

$$\rho = \bigoplus_{\alpha} p_{\alpha} \sigma_{\alpha}^{AB^L} \otimes \omega_{\alpha}^{B_{\alpha}^R C}$$

$$\Leftrightarrow \exists R_{B \rightarrow BC} \text{ s.t. } \rho_{ABC} = (\text{id}_A \otimes R_{B \rightarrow BC})(\rho_{AB})$$

called "quantum Markov state"

$$\Rightarrow \rho_{AC} \in \text{Sep}$$

Approx case?

$$\text{CMI}(\rho) \neq \min_{\sigma \text{ Markov}} D(\rho || \sigma) \quad \text{unlike classical case}$$

$$\text{but } \text{CMI} \geq -2 \log \max_R F(\rho_{ABC}, R_{B \rightarrow BC}(\rho_{AB}))$$

$$\text{i.e. } \exists R_{B \rightarrow BC} \text{ s.t. } F(\rho_{ABC}, R(\rho_{AB}))^2 \geq \exp(-\text{CMI})$$

[Fawzi - Renner '14]

more generally

$$D(\rho || \sigma) - D(N(\rho) || N(\sigma)) \geq -2 \log \max_{R \text{ s.t. } R(N(\sigma)) = \sigma} F(\rho, R(N(\rho)))$$

PF sketch

q state redistribution needs  $n I(A; C | B)$  qubits (and some cbits) for  $AC-B \Rightarrow ACB$   
 replace communication with max. mixed states. Use dF reduction to get single-letter protocol with rel. err.  $I(A; C | B)$

# Many-body physics and q. information

$H$  acts on  $(\mathbb{C}^d)^{\otimes n}$

$H = \sum_{S \subseteq [n]} h_S$   $h_S$  acts on qubits in  $S$ , eg  $H = \sum_{i \sim j} h_{ij}$  or  $\sum_{ijk} h_{ijk}$  etc..

Bosons  $d \rightarrow \infty$  Fermions later

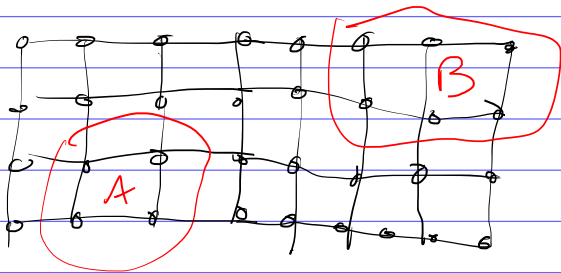
objects of study:  $e^{-iHt}$ ,  $e^{-H/T}$ ,  $| \psi_{gs} \rangle$

complexity? BQP  $\approx$  QMA

but these are worst case - often things are easier

big question How does locality of  $H$  affect locality of  $e^{-iHt}$ ,  $e^{-H/T} : | \psi_{gs} \rangle$ ?

one answer: two-point correlations



operators  $X_A, Y_B$   
 $A = \text{supp } X$  etc.

$$X_A(t) = e^{iHt} X_A e^{-iHt}$$

correlations

$$\langle X_A Y_B \rangle - \langle X_A \rangle \langle Y_B \rangle$$

another answer: full description of the state  $\cong$  general observables  
e.g. consider the toric code